

# Cluster Algebras and $q$ -Painlevé Equations

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Web-seminar on Painlevé equations and related topics

This talk is based on the collaborative research:

[1] T. Masuda, N. Okubo & TT: Birational Weyl group actions and  $q$ -Painlevé equations via mutation combinatorics in cluster algebras, arXiv:2303.06704.

[2] T. Masuda, Y. Mizuno, N. Okubo, Y. Terashima & TT: —, Part II:  $\tau$ -function formalism (a tentative title), in preparation.

# §1 Background and Results

a discrete Painlevé equation

$\cong$

birational action of a translation  
part of an affine Weyl group

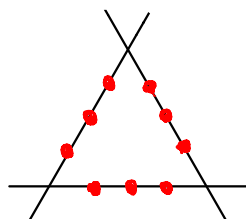
## Weyl group action arises in an algebro-geometric setup

Cf.) Dolgachev–Ortland: *Point sets in projective spaces and theta functions*. (1988)

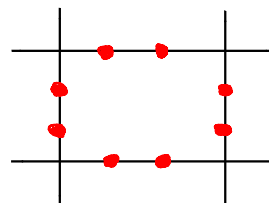
- 9 point blow-up of  $\mathbb{P}^2 \curvearrowright W(E_8^{(1)})$

$\rightsquigarrow$  Classification of 2-D Painlevé systems with the elliptic-difference Painlevé equation of type  $E_8^{(1)}$  at the top (Sakai's list).

Ref.) Sakai: Rational surfaces associated with affine root systems and geometry of the Painlevé equations. *Comm. Math. Phys.* **220** (2001)



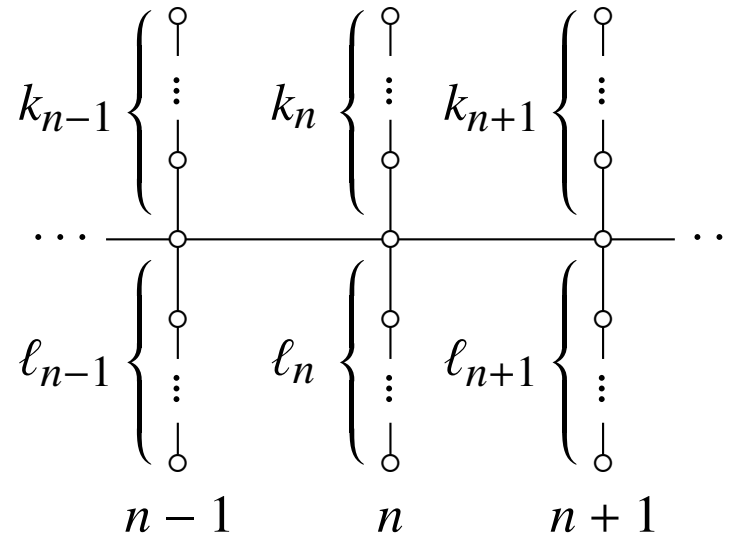
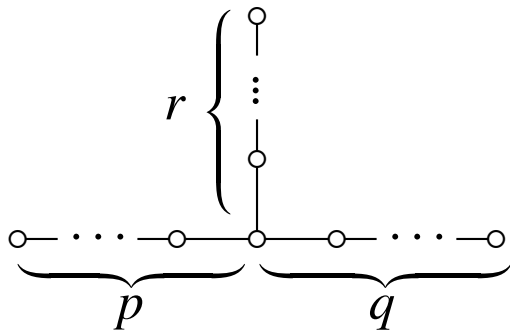
$\curvearrowright W(E_6^{(1)})$



$\curvearrowright W(D_5^{(1)})$

***q*-difference**

- It can be  $\left\{ \begin{array}{l} \text{not affine} \\ \text{higher dimensional} \end{array} \right\} \rightsquigarrow$  Various Dynkin diagrams



Ref.) Tsuda: Tropical Weyl group action via point configurations and  $\tau$ -functions of the  $q$ -Painlevé equations. *Lett. Math. Phys.* **77** (2006);

Tsuda–Takenawa: Tropical representation of Weyl groups associated with certain rational varieties. *Adv. Math.* **221** (2009)

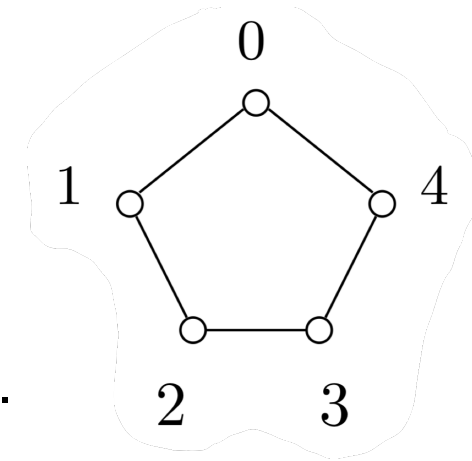
Thanks to the geometric background (= the space of initial values), these Weyl group actions admit  $\tau$ -function formalism.

E.g.)  $A_4^{(1)}$  type ( $q$ - $P_V$ )

$$\mathbb{C}(a_0, a_1, \dots, a_4; f, g) \curvearrowright W(A_4^{(1)}) = \langle s_i \ (0 \leq i \leq 4) \rangle$$

$$s_i(a_j) = a_j a_i^{-c_{ij}}, \quad s_0(f) = \frac{f}{a_0} \cdot \frac{a_2 a_3 f + a_0 g + a_0 a_3}{a_2 a_3 f + a_0^2 g + a_0 a_3}$$

etc.



$$\mathbb{C}(a_0, a_1, \dots, a_4; \tau_1, \tau_2, \dots, \tau_7) \curvearrowright W(A_4^{(1)})$$

$$(f, g) = \left( \frac{\tau_1 \tau_2}{\tau_3 \tau_7}, \frac{\tau_4 \tau_5}{\tau_6 \tau_7} \right)$$

$$s_1(\tau_{\{1,2\}}) = \tau_{\{2,1\}}, \quad s_4(\tau_{\{4,5\}}) = \tau_{\{5,4\}},$$

$$s_2(\tau_1) = \frac{a_0 a_1}{a_3^2} \frac{a_0 \tau_4 \tau_5 + a_2 a_3 \tau_6 \tau_7}{\tau_3}, \quad s_2(\tau_3) = \frac{a_0 a_1}{a_2 a_3^2} \frac{a_0 \tau_4 \tau_5 + a_3 \tau_6 \tau_7}{\tau_1},$$

$$s_3(\tau_4) = \frac{a_2}{a_0^2 a_3 a_4} \frac{a_2 a_3 \tau_1 \tau_2 + a_0 \tau_3 \tau_7}{\tau_6}, \quad s_3(\tau_6) = \frac{a_2 a_3}{a_0^2 a_4} \frac{a_2 \tau_1 \tau_2 + a_0 \tau_3 \tau_7}{\tau_4},$$

$$s_0(\tau_1) = \frac{a_4}{a_0^2 a_1 a_2} \frac{a_2 a_3 \tau_1 \tau_2 \tau_6 + a_0 \tau_3 \tau_4 \tau_5 + a_0 a_3 \tau_3 \tau_6 \tau_7}{\tau_4 \tau_7},$$

$$s_0(\tau_4) = \frac{a_0 a_4}{a_1 a_2} \frac{a_2 a_3 \tau_1 \tau_2 \tau_6 + a_0 \tau_3 \tau_4 \tau_5 + a_3 \tau_3 \tau_6 \tau_7}{\tau_1 \tau_7},$$

$$s_0(\tau_7) = \frac{a_4}{a_0 a_1 a_2} \frac{a_2 a_3 \tau_1 \tau_2 \tau_6 + a_0^2 \tau_3 \tau_4 \tau_5 + a_0 a_3 \tau_3 \tau_6 \tau_7}{\tau_1 \tau_4},$$

Laurent phenomenon and positivity, **just like cluster algebras!**

## Prior research ( $q$ -Painlevé from cluster algebras)

- Okubo and Hone–Inoue derived some individual equations.  
**d-KP, KdV**      **Y-system, T-system**

Ref.) Okubo: Discrete integrable systems and cluster algebras. RIMS Kokyuroku Bessatsu **B41** (2013);

Okubo: Bilinear equations and  $q$ -discrete Painlevé equations satisfied by variables and coefficients in cluster algebras. J. Phys. A **48** (2015);

Hone–Inoue: Discrete Painlevé equations from Y-systems. J. Phys. A **47** (2014)

- Bershtein–Gavrylenko–Marshakov derived all the 2-D  $q$ -Painlevé systems in Sakai’s list at the level of affine Weyl group symmetries.

Ref.) Bershtein–Gavrylenko–Marshakov: Cluster integrable systems,  $q$ -Painlevé equations and their quantization. J. High Energy Phys. **2018** (2018)

- Mizuno gave a geometric interpretation and complement to [BGM]’s result through toric surfaces.

Ref.) Mizuno:  $q$ -Painlevé equations on cluster Poisson varieties via toric geometry. Selecta Math. **30** (2024)

## Summary of results

[1] We present a systematic way to derive a birational representation of Weyl groups from cluster algebras, by means of a graph-combinatorial approach. Our framework covers a broad class of Dynkin diagrams and in affine case leads to (higher-order)  $q$ -Painlevé equations. The key ingredient is a mutation sequence, called the **reflection associated with a cycle graph**. Symplectic structure is also discussed.

[2] To establish  **$\tau$ -function formalism** for the above representation of Weyl groups, another new idea was needed; i.e. a **non-normalized cluster algebra** with two series of variables. Relations in Weyl groups are proven using the ‘separation formula’ and ‘synchronicity theorem’, which are fundamental tools in the theory of cluster algebras.

## §2 Cluster Mutations

Ref.) Fomin–Zelevinsky: Cluster algebras. IV. Coefficients. Compos. Math. **143** (2007)

$Q = (V, E)$  : a quiver (a directed graph)

$V = V(Q) = \{1, 2, \dots, N\}$  : the vertex set of  $Q$

$E = E(Q) \subset V \times V$  : the edge set of  $Q$

Assume that  $Q$  has no loops  $i \rightarrow i$  nor 2-cycles  $i \rightarrow j \rightarrow i$ .

$\mathbb{Q}(y_1, y_2, \dots, y_N)$  : a field of rational functions of  $\mathbf{y} = (y_1, y_2, \dots, y_N)$

## Mutation $\mu_k : (\mathbf{y}, Q) \mapsto (\mathbf{y}', Q')$ in direction $k \in V$

Mutated quiver  $Q' = \mu_k(Q)$  is defined by the procedure:

- (1) **Add** a new edge  $i \rightarrow j$  for each subgraph  $i \rightarrow k \rightarrow j$
- (2) **Reverse** the orientation of all edges including vertex  $k$
- (3) **Erase** the 2-cycles appeared.

Mutated  $y$ -variables  $\mathbf{y}' = \mu_k(\mathbf{y})$  are defined by the birational transformations:

$$y'_i = \begin{cases} y_k^{-1} & (i = k) \\ y_i y_k^{[b_{ki}]_+} (1 + y_k)^{-b_{ki}} & (i \neq k) \end{cases}$$

where  $B = (b_{ij})_{i,j=1}^N$  is the (signed) adjacency matrix of  $Q$  and

$[a]_+ = \max\{a, 0\}$  for  $a \in \mathbb{R}$ .



In general, it holds that

$$\mu_k^2 = \text{id} \quad \text{and} \quad \mu_i \circ \mu_j = \mu_j \circ \mu_i \quad \text{if} \quad b_{ij} = 0.$$

A symmetric group  $\mathfrak{S}_V$  acts on  $y$ -variables as permutations of labels:

$$\sigma(y_i) = y_{\sigma^{-1}(i)}, \quad \sigma \in \mathfrak{S}_V$$

and it holds that  $(i, j) \circ \mu_i = \mu_j \circ (i, j)$  holds.

Let  $G_Q$  be the set of compositions of mutations and permutations that keep  $Q = (V, E)$  invariant. Then  $G_Q$  provides a nontrivial group of birational transformations on  $\mathbb{Q}(y_1, y_2, \dots, y_N)$ . This is what we are interested in.



$$R_C \stackrel{\text{def}}{=} M^{-1} \circ (n-1, n) \circ M$$

## Reflection associated with a cycle graph $C$

The trident graph  $M(C)$  is invariant under a permutation  $(n-1, n)$ , thus

$$R_C(C) = M^{-1} \circ (n-1, n) \circ M(C) = M^{-1} \circ M(C) = C.$$

It is immediate from  $\mu_k^2 = (n-1, n)^2 = \text{id}$  that  $R_C^2 = \text{id}$ .

- Note that the mutation sequence  $R_C$  itself has already appeared in several areas of mathematics, including higher Teichmüller theory:

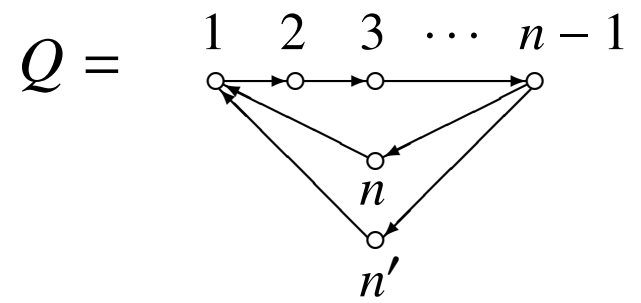
Ref.) Goncharov–Shen: Donaldson-Thomas transformations of moduli spaces of  $G$ -local systems. *Adv. Math.* **327** (2018);

Inoue–Ishibashi–Oya: Cluster realization of Weyl groups and higher Teichmüller theory. *Selecta Math.* **27** (2019);

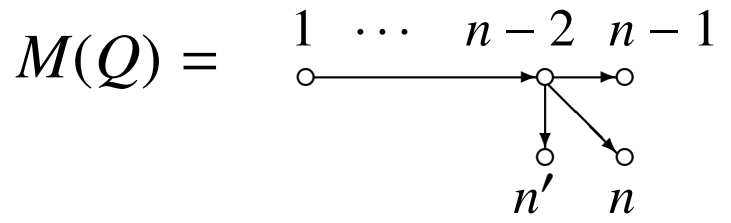
Inoue–Lam–Pylyavskyy: On the cluster nature and quantization of geometric  $R$ -matrices. *Publ. RIMS* **55** (2019);

Bucher: Maximal green sequences for cluster algebras associated to the  $n$ -torus.  
arXiv:1412.3713

Next we consider a quiver  $Q$  obtained from  $C$  by adding a copy  $n'$  of vertex  $n$ .



Applying  $M = \mu_{n-1} \circ \mu_{n-2} \circ \cdots \circ \mu_2 \circ \mu_1$ , we have



Three vertices  $n-1, n, n'$  are symmetric in  $M(Q)$ , thus  $R_C(Q) = Q$ .

**Proposition 3.1**  $(R_C \circ (n, n'))^3 = \text{id}$

**Proof**  $M$  and  $(n, n')$  are commutative. Hence we have  $R_C \circ (n, n') = M^{-1} \circ (n-1, n) \circ M \circ (n, n') = M^{-1} \circ \underbrace{(n-1, n) \circ (n, n')} \circ M$ . ■

a cyclic permutation of order 3

## Characterization of a quiver invariant under the reflections

A subgraph  $H \subseteq Q$  is called an **induced subgraph** if  $E(H)$  consists of all the edges of  $Q$  whose endpoints both belong to  $V(H)$ .

Suppose a quiver  $Q$  contains a cycle  $C = (1 \rightarrow 2 \rightarrow \cdots \rightarrow n \rightarrow 1)$  as an induced subgraph.

**Theorem 3.2**  $Q$  is  $R_C$ -invariant  $\iff Q$  satisfies the condition

(W) : For any vertex  $v$  outside of  $C$ ,  
 $\#\{\text{edges from } v \text{ to } C\} = \#\{\text{edges from } C \text{ to } v\}$

• (W) means that: any vertex  $v$  outside of  $C$  connects to  $C$  with a ‘**wedge graph**’  $w = (k \rightarrow v \rightarrow \ell)$ , where  $\{k, \ell\} \subset V(C)$ .

Cf.) [Goncharov–Shen: Donaldson-Thomas transformations of moduli spaces of  \$G\$ -local systems. Adv. Math. \*\*327\*\* \(2018\)](#) ← Sufficiency of (W) was also proven.

**Definition 3.3** A cycle subgraph  $C \subseteq Q$  is said to be **balanced** if  $C$  is an induced subgraph and satisfies the condition (W).

Hereafter, we always assume  $C \subseteq Q$  is balanced when we consider a reflection  $R_C$ .

$$\therefore R_C(Q) = Q$$

## An explicit formula of $R_C \curvearrowright \mathbb{Q}(\mathbf{y})$

$C = (1 \rightarrow 2 \rightarrow \cdots \rightarrow n \rightarrow 1)$  : an  $n$ -cycle

$$R_C = M^{-1} \circ (n-1, n) \circ M, \quad M = \mu_{n-1} \circ \mu_{n-2} \circ \cdots \circ \mu_2 \circ \mu_1$$

$$F_k(y_1, y_2, \dots, y_n) \stackrel{\text{def}}{=} 1 + y_{k+1} + y_{k+1}y_{k+2} + \cdots + y_{k+1}y_{k+2} \cdots y_{k+n-1}$$

The birational action of  $R_C$  is described as

$$R_C(y_k) = \frac{F_{k-1}}{y_{k+1}F_{k+1}} \quad (k = 1, 2, \dots, n)$$

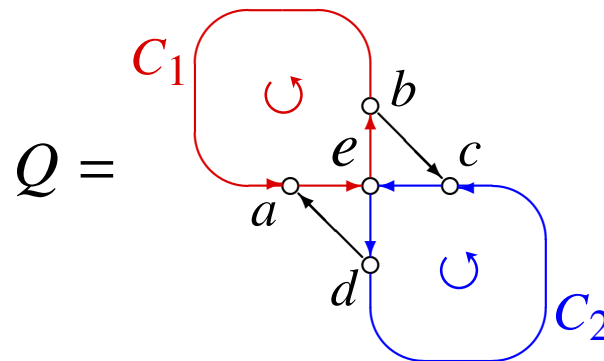
where  $y_{i+n} = y_i$ .

$$\therefore \rho^{-1} \circ R_C \circ \rho = R_C, \quad \rho = (1, 2, \dots, n) \in \mathfrak{S}_n \quad \text{Rotational symmetry}$$

$$(\curvearrowright \sigma^{-1} \circ R_C \circ \sigma = R_C, \quad \forall \sigma \in \mathfrak{S}_n \quad \text{Permutation symmetry})$$

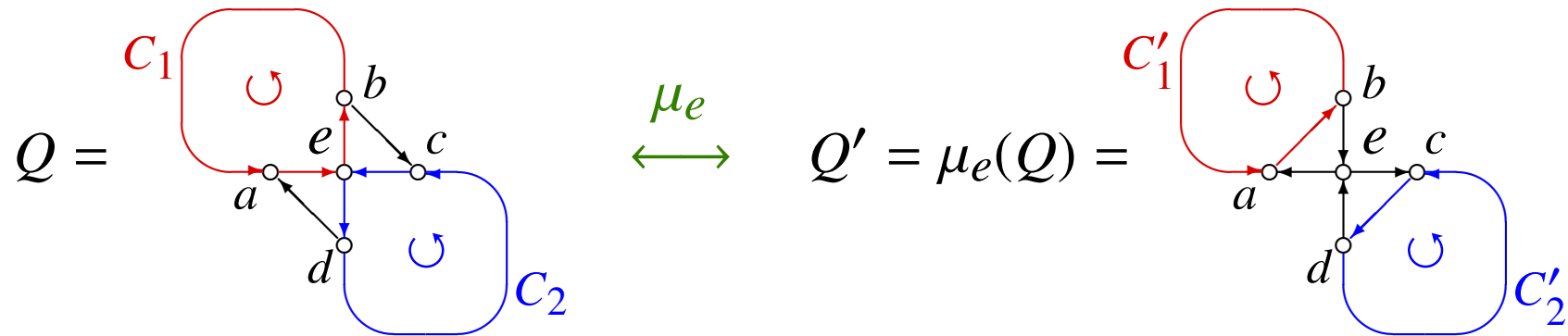
## §4 Relations of Reflections

We consider a quiver  $Q$  such that two balanced cycles  $C_1$  and  $C_2$  intersect at a vertex  $e$  as follows:



The existence of two edges  $b \rightarrow c$  and  $d \rightarrow a$  guarantees that  $C_i$  ( $i = 1, 2$ ) are balanced, thus  $R_{C_i}(Q) = Q$ .

Apply the mutation  $\mu_e$  at the crossing vertex  $e$  to 'separate'  $C_1$  and  $C_2$ .



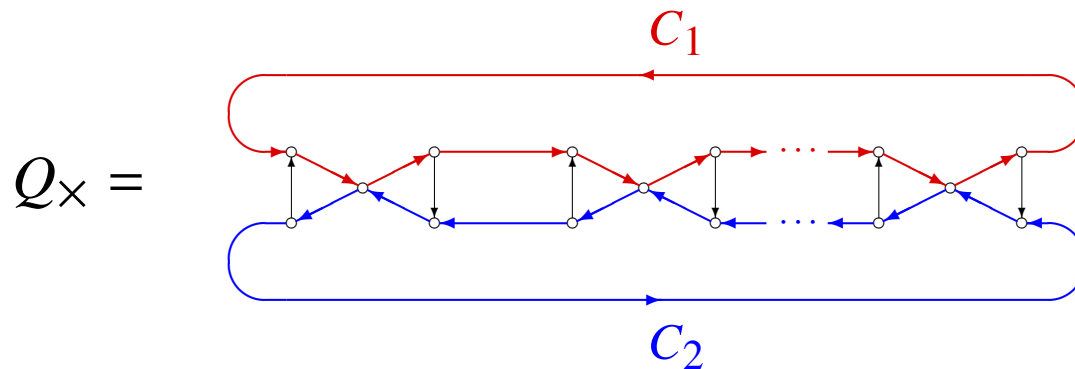
Thanks to the rotational symmetry of reflections, it holds that

$$R_{C_i} = \mu_e \circ R_{C'_i} \circ \mu_e \quad (i = 1, 2)$$

Because two cycles  $C'_1$  and  $C'_2$  are not adjacent, the reflections  $R_{C'_1}$  and  $R_{C'_2}$  mutually commute. Therefore, so do  $R_{C_1}$  and  $R_{C_2}$ :

$$\begin{aligned} R_{C_1} \circ R_{C_2} &= \mu_e \circ R_{C'_1} \circ R_{C'_2} \circ \mu_e \\ &= \mu_e \circ R_{C'_2} \circ R_{C'_1} \circ \mu_e = R_{C_2} \circ R_{C_1} \end{aligned}$$

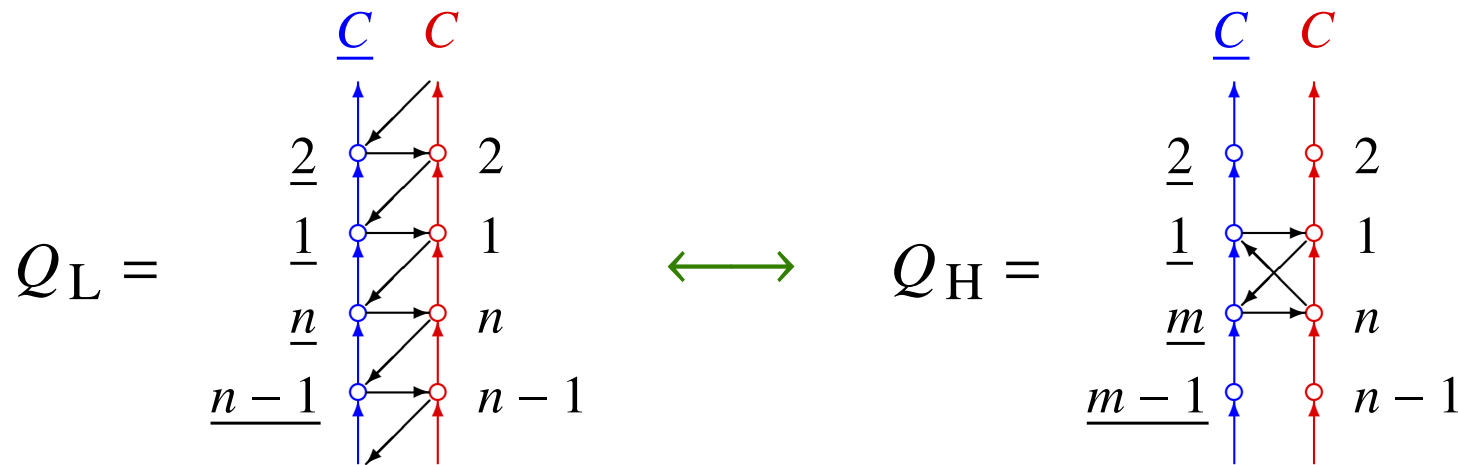
No matter how many times the two cycles  $C_1$  and  $C_2$  intersect, if we separate them by mutation at each crossing vertex, we can again deduce the commutativity of  $R_{C_1}$  and  $R_{C_2}$  from the rotational symmetry.



In general, for a quiver  $Q \supseteq Q_x$  (as an induced subgraph) we have

**Theorem 4.1** Reflections associated with two intersecting cycles are commutative; i.e.  $R_{C_1}R_{C_2} = R_{C_2}R_{C_1}$  if  $C_1 \cap C_2 \neq \emptyset$ .

- Let  $Q_L$  be a quiver consisting of two  $n$ -cycles  $\underline{C}$  and  $C$  connected with consecutive wedge graphs  $\underline{i} \rightarrow i \rightarrow \underline{i-1}$  ( $i \in \mathbb{Z}/n\mathbb{Z}$ ) like a **ladder**.
- Let  $Q_H$  be a quiver consisting of  $m$ -cycle  $\underline{C}$  and  $n$ -cycle  $C$  connected with a 'hinge'  $1 \rightarrow \underline{m} \rightarrow n \rightarrow \underline{1} \rightarrow 1$ .



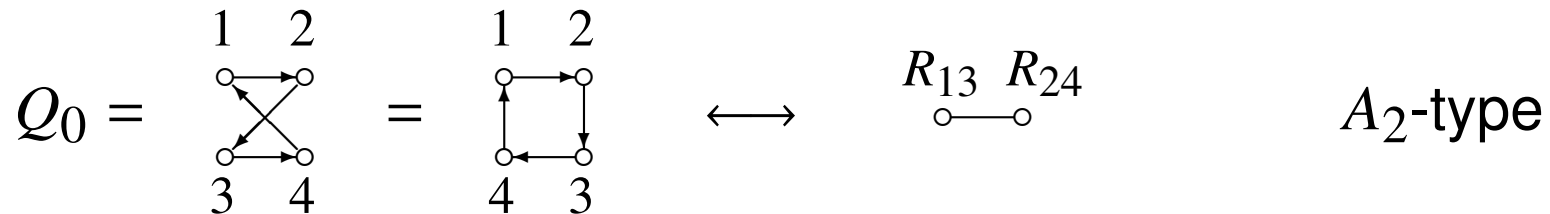
**Theorem 4.2** In both cases (i)  $Q \supseteq Q_L$  and (ii)  $Q \supseteq Q_H$ , the braid relation  $R_C R_{\underline{C}} R_C = R_{\underline{C}} R_C R_{\underline{C}}$  ( $\Leftrightarrow (R_C R_{\underline{C}})^3 = \text{id}$ ) holds.

**Proof** Due to the rotational symmetry, both cases reduce to a much simpler result via mutations: (i)  $\rightsquigarrow$  (ii)  $\rightsquigarrow$  Prop 3.1 (= **Cycle & a copy**) ■

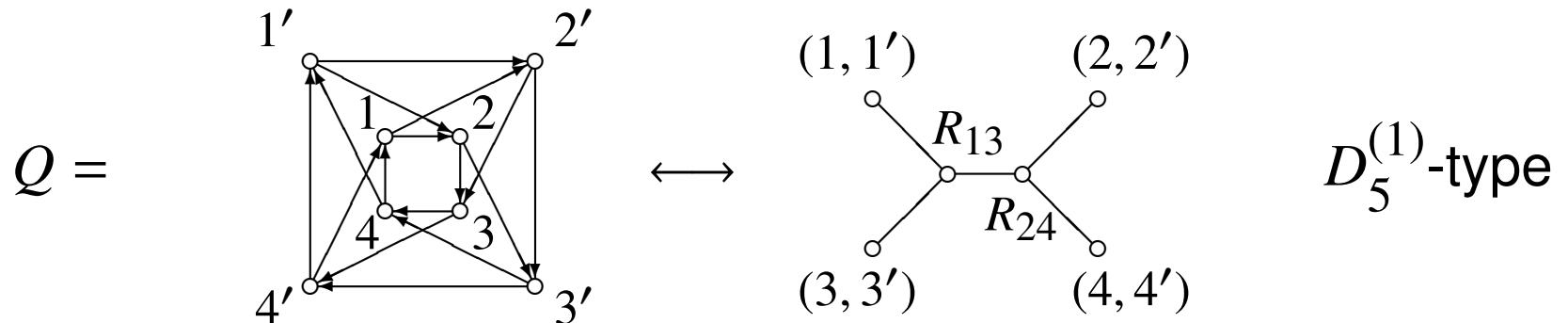
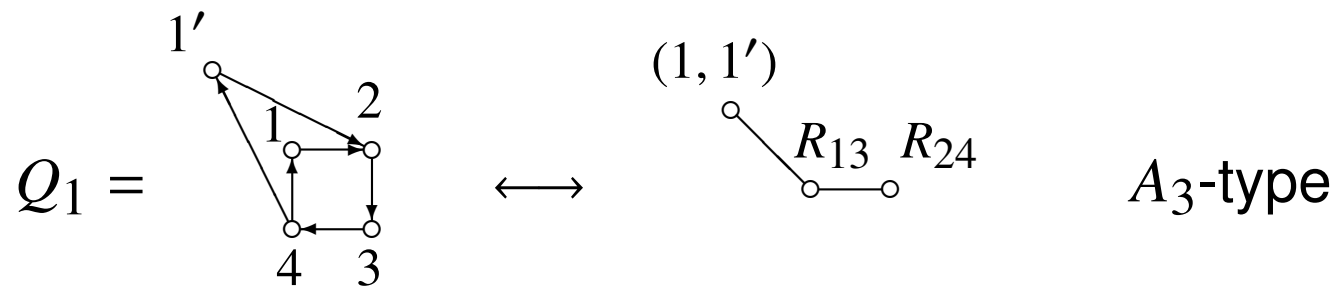
# §5 Examples

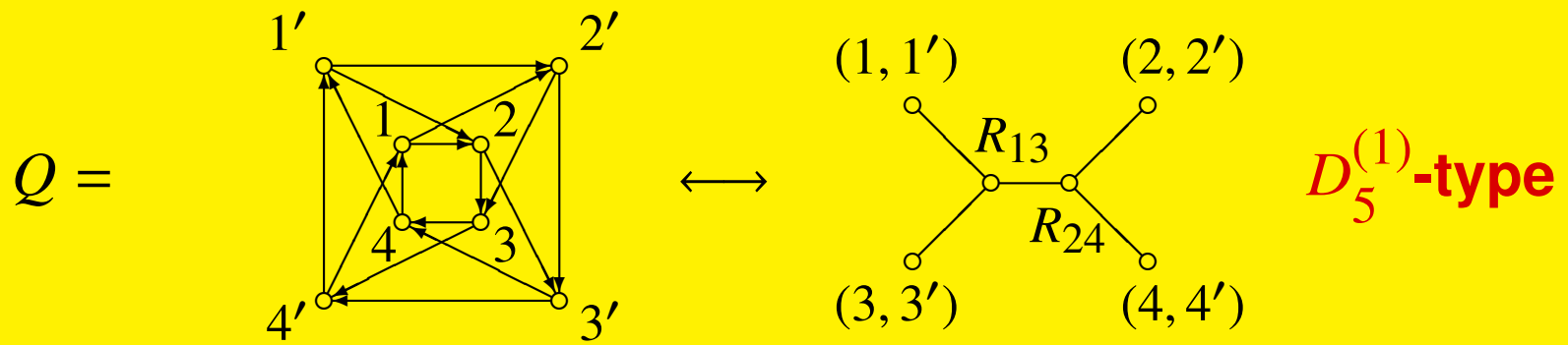
$q$ - $P_{VI}$

Adjacent 2-cycles  $C_{13} = (1 \rightarrow 3 \rightarrow 1)$  and  $C_{24} = (2 \rightarrow 4 \rightarrow 2)$ :



Add a copy  $i'$  of vertex  $i$  to  $Q_0$ :





Birational action of a translation in the affine Weyl group

$$W(D_5^{(1)}) = \langle R_{13}, R_{24}, \underbrace{(1, 1'), (2, 2'), (3, 3'), (4, 4')}_{\text{transpositions of vertices}} \rangle \subset G_Q$$

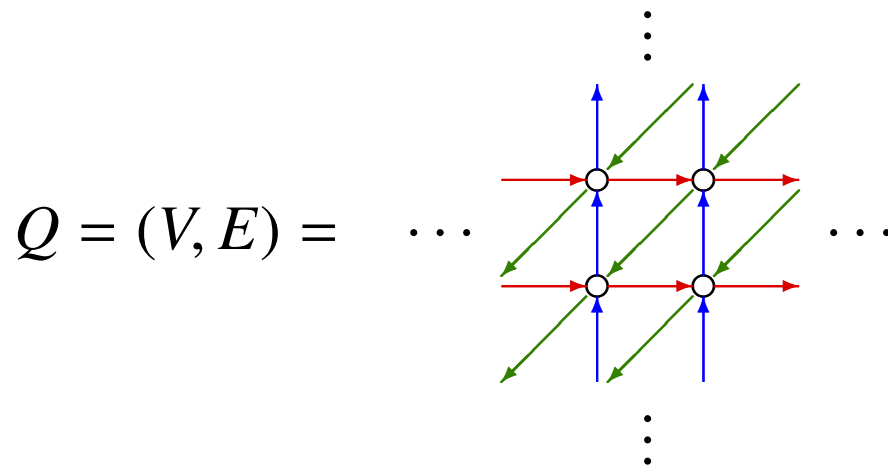
coincides with  $q$ - $P_{VI}$ .

- In a similar manner, we can construct birational realizations of Weyl groups corresponding to various Dynkin diagrams.

Higher dimensional case  $m, n \geq 2$  (and one of them  $> 2$ )

$$V = \{v_{i,j} \mid i \in \mathbb{Z}/m\mathbb{Z}, j \in \mathbb{Z}/n\mathbb{Z}\}$$

$$E = \{v_{i,j} \rightarrow v_{i+1,j}, v_{i,j} \rightarrow v_{i,j+1}, v_{i+1,j+1} \rightarrow v_{i,j}\}$$



The quiver  $Q$  contains the following balanced cycles:

- **Vertical**  $C_i^v = (v_{i,1} \rightarrow v_{i,2} \rightarrow \dots \rightarrow v_{i,n} \rightarrow v_{i,1}) \quad (i \in \mathbb{Z}/m\mathbb{Z})$
- **Horizontal**  $C_j^h = (v_{1,j} \rightarrow v_{2,j} \rightarrow \dots \rightarrow v_{m,j} \rightarrow v_{1,j}) \quad (j \in \mathbb{Z}/n\mathbb{Z})$

If  $m$  and  $n$  are not relatively prime, there is another type ( $g = \gcd(m, n)$ ):

- **Diagonal**  $C_k^d = (v_{k,0} \rightarrow v_{k-1,-1} \rightarrow v_{k-2,-2} \rightarrow \dots \rightarrow v_{k,0}) \quad (k \in \mathbb{Z}/g\mathbb{Z})$

Reflections  $s_i^\Delta = R_{C_i^\Delta}$  ( $\Delta = v, h, d$ ) keeps  $Q$  invariant.

- Theorem 4.1 implies that  $s_i^\Delta s_j^{\Delta'} = s_j^{\Delta'} s_i^\Delta$  ( $\forall i, j$ ) if  $\Delta \neq \Delta'$ .
- Theorem 4.2 implies that  $(s_i^\Delta s_{i+1}^\Delta)^3 = \text{id}$  for each  $\Delta = v, h, d$ .

$$\therefore W(A_{m-1}^{(1)}) \times W(A_{n-1}^{(1)}) \times W(A_{g-1}^{(1)}) \simeq \langle s_i^v, s_i^h, s_i^d \rangle \curvearrowright \mathbb{Q}(\{y_{i,j}\})$$

**E.g.)** Reflections associated with '**vertical cycles**' are described as

$$s_i^v(y_{i,j}) = \frac{F_{i,j-1}}{y_{i,j+1}F_{i,j+1}}, \quad s_i^v(y_{i+1,j}) = \begin{cases} \frac{s_i^v(y_{i-1,j-1})}{y_{i-1,j-1}} = \frac{y_{i,j}F_{i,j}}{F_{i,j-1}} & (\text{if } m \geq 3) \\ \frac{y_{i,j}y_{i,j+1}F_{i,j+1}}{F_{i,j-1}} & (\text{if } m = 2) \end{cases}$$

where  $F_{i,j} = 1 + \sum_{a=1}^{n-1} \prod_{b=1}^a y_{i,j+b}$ . (**Horizontal** and **diagonal** are similar.)

- The **vertical** and **horizontal** parts, i.e.  $W(A_{m-1}^{(1)}) \times W(A_{n-1}^{(1)})$ , are equivalent to those realized by Yamada and Kajiwara–Noumi–Yamada.

Ref.) Yamada: A birational representation of Weyl group, combinatorial  $R$ -matrix and discrete Toda equation. (in Proceedings of Physics and Combinatorics 2000);  
 Kajiwara–Noumi–Yamada: Discrete dynamical systems with  $W(A_{m-1}^{(1)} \times A_{n-1}^{(1)})$  symmetry. Lett. Math. Phys. **60** (2002)

- Based on our framework including **diagonal**, Okubo–Suzuki showed that various higher-order  $q$ - $P_{\text{VI}}$ , such as the  $q$ -Garner system, can be derived from the same  $Q$  if  $m = 2$  and  $n$  is even greater than two.

Ref.) Okubo–Suzuki: Generalized  $q$ -Painlevé VI systems of type  $(A_{2n+1} + A_1 + A_1)^{(1)}$  arising from cluster algebra. IMRN **2022** (2022)

- In fact, the symmetry of  $Q$  increases only in the exceptional cases:

$$(m, n) = (2, 4), (4, 2) \Rightarrow D_5^{(1)}, \quad (m, n) = (3, 3) \Rightarrow E_6^{(1)}$$

## §6 Reduction of Dynamical Systems

$B = (b_{ij})_{i,j=1}^N$  : the adjacency matrix of  $Q$   
 (= a skew symmetric integer matrix)

$$\{y_i, y_j\} \stackrel{\text{def}}{=} b_{ij} y_i y_j \quad \text{Poisson bracket}$$

Ref.) Gekhtman–Shapiro–Vainshtein: *Cluster Algebras and Poisson Geometry*. (2010)

With a multi-index notation  $y^{\mathbf{m}} = \prod_{i=1}^N y_i^{m_i}$ , Leibniz's rule shows

$$\{y^{\mathbf{m}}, y^{\mathbf{n}}\} = (\mathbf{m}^T B \mathbf{n}) y^{\mathbf{m}+\mathbf{n}}.$$

$\therefore y^{\mathbf{m}}$  is a Casimir function  $\iff \mathbf{m} \in \ker B$

**Lemma 6.1** If  $\mathbf{m} \in \ker B \cap \mathbb{Z}^N$  then

$$\mu_k(y^{\mathbf{m}}) = y^{A_k \mathbf{m}} \quad \text{and} \quad A_k \mathbf{m} \in \ker \mu_k(B) \cap \mathbb{Z}^N$$

**Proof** The mutation rule  $\mu_k(B) = {}^T A_k B A_k$ ,  $\mu_k(y^{\mathbf{m}}) = y^{A_k \mathbf{m}} (1+y_k)^{-(B\mathbf{m})_k}$  and  $A_k^2 = \text{id}$  lead us to the desired result. ■

## Lemma 6.2 (Normal form of a skew-symmetric integer matrix $B$ )

$\exists U = (u_{ij})_{i,j=1}^N \in GL_N(\mathbb{Z})$  s.t.

$${}^T U B U = \begin{pmatrix} 0 & h_1 \\ -h_1 & 0 \end{pmatrix} \oplus \begin{pmatrix} 0 & h_2 \\ -h_2 & 0 \end{pmatrix} \oplus \cdots \oplus \begin{pmatrix} 0 & h_\ell \\ -h_\ell & 0 \end{pmatrix} \oplus O_{N-2\ell},$$

where the positive integers  $h_1, h_2, \dots, h_\ell$  satisfying  $h_i | h_{i+1}$  are unique.

Ref.) Newman: *Integral Matrices*. (1972)

Write  $U$  by arranging column vectors  $\mathbf{u}_i$  ( $1 \leq i \leq N$ ) in a row.

$$U = (\mathbf{u}_1, \dots, \mathbf{u}_{2\ell}, \underbrace{\mathbf{u}_{2\ell+1}, \dots, \mathbf{u}_N}_{\ker B}) \in GL_N(\mathbb{Z}).$$

- The first  $2\ell$  half  $\rightsquigarrow f_i = y^{u_{2i-1}}, g_i = y^{u_{2i}}$  ( $1 \leq i \leq \ell$ ) **‘Canonical coordinates’**
- The latter  $N-2\ell$  half ( **$\ker B$** )  $\rightsquigarrow \kappa_i = y^{u_{2\ell+i}}$  ( $1 \leq i \leq N-2\ell$ ) **Casimir**

The new variables  $\{f_i, g_i, \kappa_i\}$  are Laurent monomials in the original  $\{y_i\}$ , and vice versa ( $\because$  unimodularity of  $U$ ).

Lemma 6.2 implies

$$\{f_i, f_j\} = \{g_i, g_j\} = 0, \quad \{f_i, g_j\} = h_i \delta_{ij} f_i g_j$$

‘Canonical commutation relations’

A transformation  $\varphi$  of  $(f_i, g_i)$  preserves  $\{\cdot, \cdot\}$

$\iff$

$\varphi$  preserves the 2-form

$$\omega = \sum_{i=1}^{\ell} \frac{1}{h_i} \frac{df_i \wedge dg_i}{f_i g_i}$$

In general, it holds that  $\{\mu_k(y_i), \mu_k(y_j)\} = \mu_k(b_{ij}) \mu_k(y_i) \mu_k(y_j)$ . Hence, in particular,  $\forall \varphi \in G_Q$  preserves the Poisson bracket  $\{\cdot, \cdot\}$ .

**Theorem 6.3**  $G_Q$  gives a group of birational canonical transformations on  $\mathbb{K}(f_1, \dots, f_\ell, g_1, \dots, g_\ell)$  in the sense that  $\forall \varphi \in G_Q$  preserves the 2-form  $\omega$ . Here the coefficient field is  $\mathbb{K} = \mathbb{Q}(\kappa_1, \dots, \kappa_{N-2\ell})$  and the number of canonical coordinates is  $2\ell = \text{rank} B$ .

$\{\eta_1, \eta_2, \dots, \eta_r\}$  : a basis of  $\ker B$ ,  $r = N - 2\ell$

**Assume:** Each  $\eta_i \in \mathbb{R}^N$  is either of the following two types:

- (i)  $\varepsilon_C = \sum_{i \in V(C)} e_i$  **cycle type**;      (ii)  $e_i - e_j$  **transposition type**

where  $C \subseteq Q$  is a balanced cycle and  $e_i = \text{T}(0, \dots, 0, \overset{i}{1}, 0, \dots, 0) \in \mathbb{R}^N$ .

Introduce the (multiplicative) **root variables**  $a_k = y^{\eta_k}$  attached to a basis  $\eta_k$  ( $1 \leq k \leq r$ ) of  $\ker B$ . Define the corresponding reflections as

$$s_k = \begin{cases} R_C & \text{if } \eta_k \text{ is of cycle type (i)} \\ (i, j) & \text{if } \eta_k \text{ is of transposition type (ii).} \end{cases}$$

Then the transformation law takes the form  $s_i(a_j) = a_j a_i^{-c_{ij}}$ .

(Furthermore, under a certain technical assumption)

**Theorem 6.4**  $A = (c_{ij})_{i,j=0}^{r-1}$  is a generalized Cartan matrix (GCM) i.e.

- (i)  $c_{ii} = 2$       (ii)  $c_{ij} \in \mathbb{Z}_{\leq 0}$  ( $i \neq j$ )      (iii)  $c_{ij} = 0 \Leftrightarrow c_{ji} = 0$

## §7 $\tau$ -Function Formalism (Outline)

$Q = (V, E)$ ,  $B = (b_{ij})_{i,j=1}^N$  : a quiver and its adjacency matrix

|           | Seed                          | Mutations                                      |
|-----------|-------------------------------|--|
| Until now | $(Q, \mathbf{y})$             | $\mu_k$  |
| From now  | $(Q, \mathbf{p}, \mathbf{x})$ | $\mu_k^{(\gamma)}$ ( $\gamma \in \mathbb{R}$ ) |

- Quiver mutation is the same as before:  $Q' = \mu_k^{(\gamma)}(Q) = \mu_k(Q)$
- Variable mutations  $\mathbf{x}' = \mu_k^{(\gamma)}(\mathbf{x})$  and  $\mathbf{p}' = \mu_k^{(\gamma)}(\mathbf{p})$  :

$$x'_i = \begin{cases} \frac{p_k^\gamma \prod_{j=1}^N x_j^{[b_{jk}]_+} + p_k^{\gamma-1} \prod_{j=1}^N x_j^{[-b_{jk}]_+}}{x_k} & (i = k) \\ x_i & (i \neq k) \end{cases}$$

$$p'_i = \begin{cases} p_k^{-1} & (i = k) \\ p_i p_k^{\gamma[b_{ki}]_+ + (1-\gamma)[b_{ik}]_+} & (i \neq k) \end{cases}$$

In general, it holds that

$$\mu_k^{(\gamma)} \circ \mu_k^{(1-\gamma)} = \text{id} \quad \text{and} \quad \mu_i^{(\gamma)} \circ \mu_j^{(\delta)} = \mu_j^{(\delta)} \circ \mu_i^{(\gamma)} \quad \text{if } b_{ij} = 0.$$

Note that  $\hat{y}_i = p_i \prod_{j=1}^N x_j^{b_{ji}}$  obey the same mutation rule as  $y_i$ .

Let  $C = (1 \rightarrow 2 \rightarrow \dots \rightarrow n \rightarrow 1) \subseteq Q$  be a balanced cycle. For any  $\boldsymbol{\gamma} = (\gamma_1, \gamma_2, \dots, \gamma_{n-1}) \in \mathbb{R}^{n-1}$ , we define

$$R_C[\boldsymbol{\gamma}] = \mu_1^{(1-\gamma_1)} \circ \dots \circ \mu_{n-1}^{(1-\gamma_{n-1})} \circ (n-1, n) \circ \mu_{n-1}^{(\gamma_{n-1})} \circ \dots \circ \mu_1^{(\gamma_1)}$$

**Reflection associated with a cycle graph  $C$**

**Lemma 7.1 (Rotational symmetry of  $R_C$  with parameter change)**

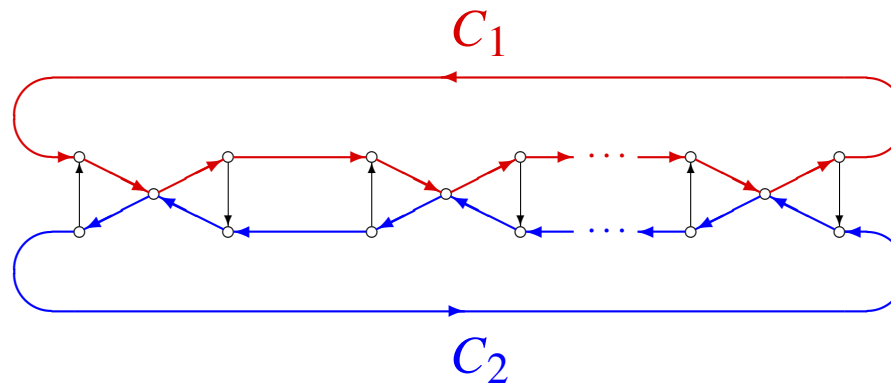
$$R_C[\gamma] = \rho^{-\ell} \circ R_C[\gamma'] \circ \rho^\ell, \quad \rho = (1, 2, \dots, n) \in \mathfrak{S}_n$$

where

$$\gamma'_i = \frac{\Gamma_{i+\ell} - \Gamma_\ell}{\Gamma_{i+\ell+1} - \Gamma_\ell},$$

with  $\Gamma_0 = 0$ ,  $\Gamma_i = \prod_{j=i}^{n-1} \gamma_j$  ( $1 \leq i \leq n-1$ ),  $\Gamma_{i+n} = \Gamma_i + 1$  ( $i \in \mathbb{Z}$ ).

As similar to §4, this rotational symmetry is a crucial key to proving the relations of reflections. For example, we have



**Theorem 7.2** Let  $C_1$  and  $C_2$  be two intersecting cycles. It holds that

$$R_{C_1}[\gamma]R_{C_2}[\delta] = R_{C_2}[\delta]R_{C_1}[\gamma]$$

if parameters  $(\gamma, \delta)$  satisfy the constraint

$$\sum_{u=v \in C_1 \cap C_2} (\Gamma_{u-1}(\Delta_{v+1} - \Delta_v) + \Gamma_u(\Delta_{v-1} - \Delta_{v+1}) + \Gamma_{u+1}(\Delta_v - \Delta_{v-1})) = 0.$$

As before,  $\Gamma_i$  and  $\Delta_i$  are certain monomials in  $\gamma$  and  $\delta$ , respectively.

**Thank you for your attention!**