## High-field magnetic properties of caged rare-earth borides $RB_n$ (n = 4, 12)

<u>F. Iga</u><sup>1</sup>, Michimura S.<sup>1</sup>, Murakami K.<sup>1</sup>, Takabatake T.<sup>1</sup>, Suga K.<sup>2</sup>, Kindo K.<sup>2</sup>, Yoshii S.<sup>3</sup>

The Rare earth borides exhibit a variety of physical properties ranging from superconductor (YB<sub>6</sub>, ZrB<sub>12</sub>), to typical Kondo lattice system with a quadrupole ordering (CeB<sub>6</sub>), valence fluctuating system with a correlated small energy gap (SmB<sub>6</sub> (valence=2.6), YbB<sub>12</sub> (valence=2.9-3.0)), multi-step magenetization system RB<sub>4</sub> (R=Tb-Tm). Single crystals of these rare earth borides have been prepared by a floating zone method using an image furnace with four xenon lamps.

YbB<sub>12</sub> is called "Kondo semiconductor" in which an energy gap of 200 K gradually opens as temperature is decreased below 100 K [1]. Application of magnetic field of approximately 50 Tesla destroys the gap and induces a first-order transition from a semiconducting state to a metallic state [2]. The magnetization of the single crystalline sample increases suddenly in an anisotropic way at the critical fields  $B_{C1} = 47$  Tesla for B//[100] and  $B_{C2} = 54$  Tesla for B//[110] and [111]. In the present study, we have measured both magnetization and magnetoresistance up to 68 Tesla by using a long-pulse magnet (pulse width 36 ms). Thereby, heating effect was much reduced. Magnetization M(B) at 1.3 K reproduced the previous data described above with hysteretic behaviors at  $B_{C1}$  and  $B_{C2}$ . Above 60 Tesla, M(B) increases linearly up to the highest field of 68 Tesla as shown in Fig. 1.

Rare-earth tetraboride  $RB_4$  has a tetragonal crystal structure where the sublattice of R ions in the c-plane is equivalent to the Shastry-Sutherland lattice, that is known as a model of a quantum spin system  $SrCu_2(BO_3)_2$  [3]. A novel magnetic ordering state is expected by the geometric frustration between the orthogonal R-R dimmers as seen in  $DyB_4$  and  $HoB_4$  [4].

We have determined magnetic phase diagrams of isostructural system  $TbB_4$ ,  $ErB_4$  and  $TmB_4$  from the measurements of magnetization, magnetoresistance, magnetostriction, specific heat, neutron diffraction and Hall effect. Magnetization process for B//c and a in  $ErB_4$  [5] and  $TmB_4$  [6] shows a plateau at half moment and the full saturated moment, respectively. Otherwise,  $TbB_4$  shows nine plateaus for B//c [7], but a simple metamagnetic transition for B//a. These field dependences and obtained magnetic phase diagrams would be explained as a result of competition between the frustration effect and Zeeman effects.

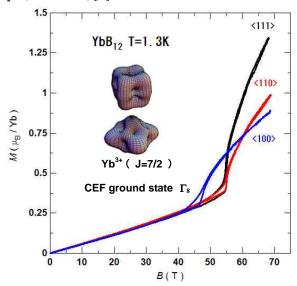


Fig. 1 Magnetization process at 1.3 K of  $YbB_{12}$  for B/[100], [110] and [111] up to 68 T.

## References

- 1) F. Iga et al., J. Mag. Mag. Mater. 177-181 (1998) 53.
- 2) S. Kawasaki et al., Physica B 281&282 (2000) 234.
- 3) H. Kageyama, et al., Phys. Rev. Lett. 82 (1999) 3168.
- 4) R. Watanuki et al., J. Phys. Soc. Jpn. 74 (2005) 2169.
- 5) S. Michimura *et al.*, Physica B **378-380** (2006), 596.
- 6) F. Iga et al., J. Magn. Magn. Mater. 310 (2007) e443.
- 7) S. Yoshii et al., J. Magn. Magn. Mater. 310 (2007) 1282.

<sup>&</sup>lt;sup>1</sup> ADSM, Hiroshima University, Higashi-Hiroshima, Hiroshima 739-8530

<sup>&</sup>lt;sup>2</sup> Institute for Solid State Physics, University of Tokyo, Kashiwa, Chiba 277-8581

<sup>&</sup>lt;sup>3</sup> Institute for Materials Research, Tohoku University, Sendai, Miyagi 980-8577