

(29b1)

## Theory of Superconductivity in Heavy Fermion Systems

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In this presentation, we show the results obtained within the third-order perturbation theory for superconductivity in some typical heavy fermion compounds. This approach has been useful in discussing the pairing symmetry and mechanism of superconductivity in high- $T_c$  cuprates and  $\text{Sr}_2\text{RuO}_4$ . Although the perturbation theory with respect to the on-site repulsion  $U$  is the approach from a weak correlation limit, we indicate here that it is valuable also for heavy fermion superconductors. It is important that the normal state above  $T_c$  is the Fermi liquid state with heavy electron mass, and the large mass enhancement factor does not have remarkable wave vector dependence as implied by the Kadowaki-Woods relation. In this case, we discuss superconductivity by starting with the renormalized quasi-particle picture. The obtained results are illustrated below.

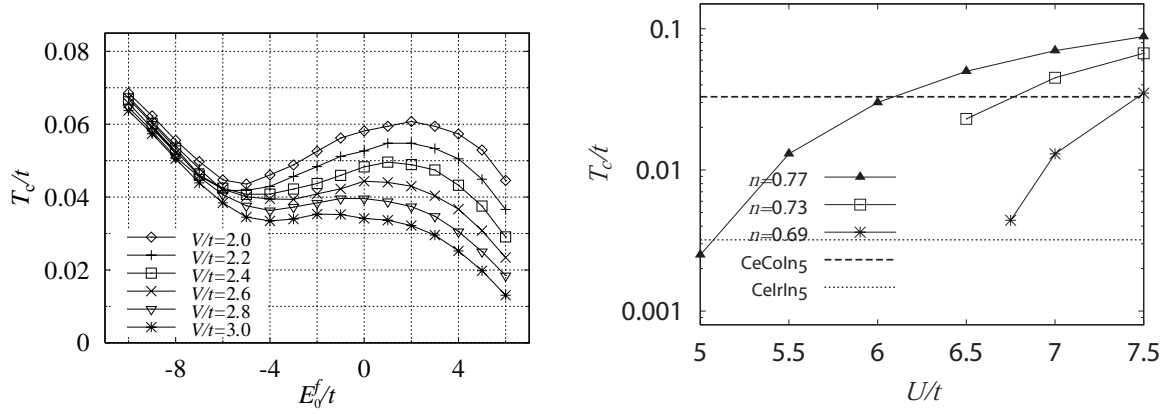


Figure 1: Left:  $f$ -level  $E_0^f$  dependence of  $T_c$  within the third-order perturbation theory in the periodic Anderson model for  $\text{CeCu}_2(\text{Si,Ge})_2$ . Right:  $U$  dependence of  $T_c$  for some electron density  $n$  in the Hubbard model on the 2D square lattice. The tendency that  $T_c$  increases with increasing  $n$  is consistent with its  $x$ -dependence in  $\text{CeIr}_x\text{Co}_{1-x}\text{In}_5$ .

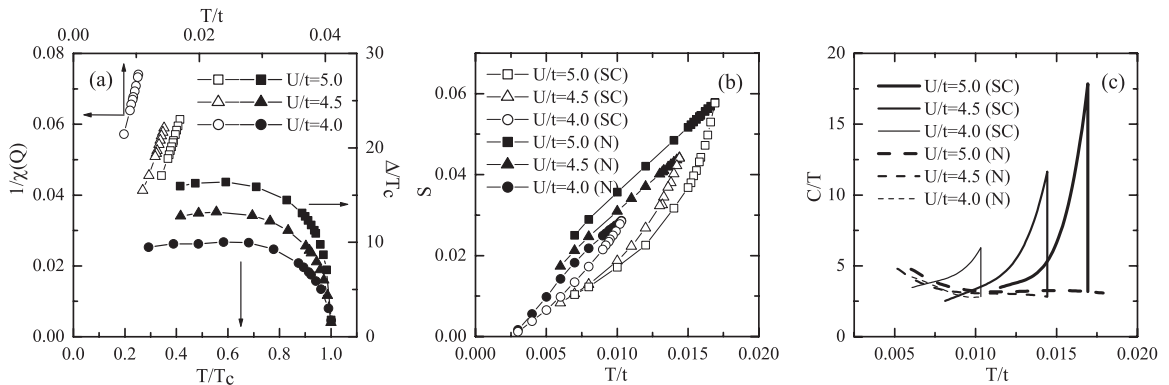


Figure 2: (a) Temperature dependence of  $1/\chi(Q)$  and the magnitude of the gap function  $\Delta/T_c$ . They are obtained within the fluctuation-exchange approximation. (b) The entropy  $S$  and (c) the specific heat  $C$  in the normal and the superconducting states. As  $U/t$  is larger, the superconductivity becomes the strong coupling, and the jump of  $C/T$  at  $T_c$  is enhanced. This is consistent with the large specific heat jump in  $\text{CeCoIn}_5$ .

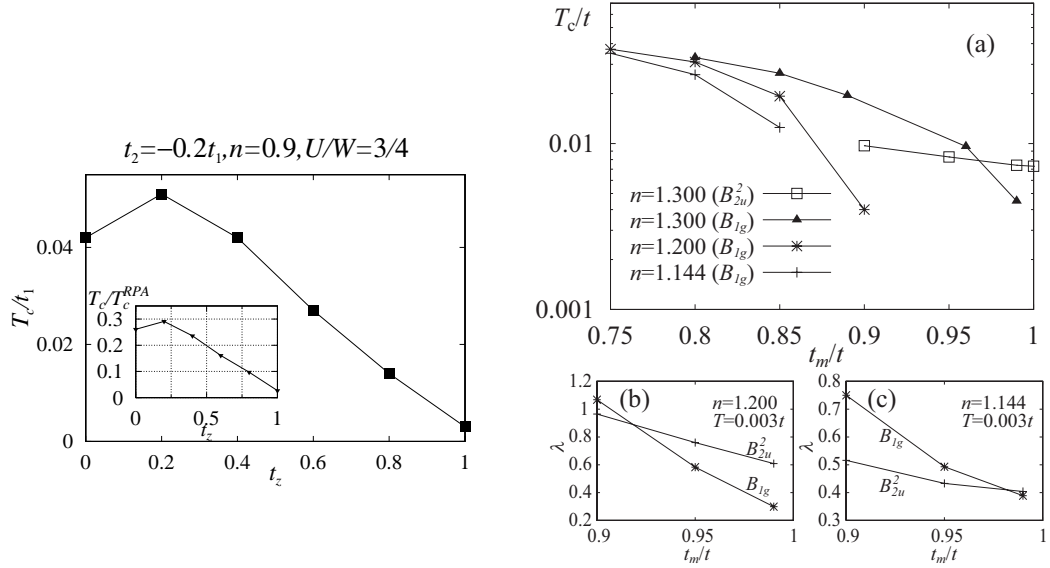


Figure 3: Left:  $t_z$  dependence of  $T_c/t_1$  in the 3D cubic Hubbard model. In the inset,  $T_c/T_c^{\text{RPA}}$ -like as a function of the anisotropy  $t_z$ .  $t_z = 0$  and 1 correspond to 2D and 3D systems, respectively.  $T_c$  in 3D systems becomes one-order smaller than that in 2D systems. This tendency corresponds to difference of  $T_c$  between  $\text{CeCoIn}_5$  and  $\text{CeIn}_3$ . Right: (a)  $T_c$  as a function of the distortion ratio  $t_m/t$  in the Hubbard model on a distorted triangular lattice. (b) and (c) illustrate, respectively, the eigen values  $\lambda$  of the Éliashberg equation for  $n = 1.200$  and  $n = 1.144$  at  $T = 0.003t$ .  $\Delta_{B_{1g}}$  ( $d_{xy}$ -wave singlet) is stable in a wide range of  $t_m/t$ , while  $\Delta_{B_{2u}^2}$  ( $p$ -wave triplet) only in the vicinity of  $t_m/t = 1$ . Although the region of  $\Delta_{B_{2u}^2}$  becomes wider as  $n$  decreases, the transition temperature is suppressed abruptly. This can correspond to the fact that singlet pairing with line nodes in  $\text{UPd}_2\text{Al}_3$  and triplet pairing in  $\text{UNi}_2\text{Al}_3$  are realized.

[1] Y. Yanase, T. Jujo, T. Nomura, H. Ikeda, T. Hotta and K. Yamada, Phys. Rep. **387** (2003) 1