

(8b4)

**f-electron states of  $\text{RT}_4\text{X}_{12}$  ( $\text{R}=\text{Ce}\sim\text{Yb}$ )**

Takashi Hotta

Advanced Science Research Center, JAERI, Tokai, Ibaraki 319-1195

In order to clarify magnetic properties of filled skutterudites, we analyze the Anderson model including seven  $f$  orbitals hybridized with an  $a_u$  conduction band using a numerical technique. For  $n=2$  corresponding to Pr-based filled skutterudites, where  $n$  is the local  $f$ -electron number, even if the ground state is a  $\Gamma_1$  singlet, there remain significant magnetic fluctuations from a  $\Gamma_4^{(2)}$  triplet state with a small excitation energy. This result can be understood by the fact that  $f$ -electron states are distinguished as itinerant  $\Gamma_7$  and localized  $\Gamma_8$  in the filled skutterudite structure, since the  $a_u$  conduction band with xyz symmetry is described by  $\Gamma_7$ . This picture also explains the complex results for  $f$ -electron magnetic susceptibility and entropy for  $n=1\sim 13$ .

In Figs. 1(a) and (b), we show the numerical results of susceptibility  $\chi$  for the cases of  $n=1\sim 13$ . Note that we always use the same CEF parameters determined for  $n=2$ . For  $n=1, 3$ , and 5, the ground state is a  $\Gamma_{67}$  quartet in  $T_h$ , as confirmed from the residual entropy of  $\log 4$ . For  $n=4$  and 6, the ground state is a  $\Gamma_1$  singlet. The excitation energy is large and both magnetic and orbital fluctuations should be rapidly suppressed with decreasing temperature. Thus,  $T\chi$  immediately becomes zero at low temperatures for  $n=4$  and 6. For  $n\geq 7$ , the absolute values of  $\chi$  are much larger than those for  $n<7$ , since the total angular momentum  $J$  becomes large for  $n\geq 7$  owing to Hund's rule coupling. Typically, at half-filling, total spin  $S(=J)$  is equal to  $7/2$ , and the Curie constant for an isolated ion is as large as  $21\mu_B^2/k_B$ . Over a broad temperature region, this value has been observed for  $n=7$ , indicating that the  $S=7/2$  spin survives at relatively low temperatures. For  $n=8$  and 12, the ground state is a  $\Gamma_1$  singlet and the magnetic excited state energy is now large, in sharp contrast with the case of  $n=2$ . Thus, the susceptibility rapidly goes to zero. For  $n=9, 11$ , and 13, the local ground state is a  $\Gamma_5$  Kramers doublet in  $T_h$ , which is the mixture of  $\Gamma_6$  and  $\Gamma_7$  of  $O_h$ . Since  $\Gamma_6$  state does not hybridize with the  $a_u$  conduction band, the magnetic moment from  $\Gamma_6$  still persists even in the low-temperature region. In fact, we observe a residual entropy of  $\log 2$  in these cases. For  $n=10$ , the local ground state is a  $\Gamma_4$  triplet, but the local triplet seems to remain at low temperatures. This is easily understood, if we recall that  $\Gamma_4$  triplets of  $T_h$  are given by mixtures of  $\Gamma_4$  and  $\Gamma_5$  of  $O_h$ . The  $\Gamma_5$  triplet still persists even after hybridization, since it is composed of a couple of  $\Gamma_8$  electrons and a  $\Gamma_8$  electron does not hybridize with a  $\Gamma_7$  conduction electron.

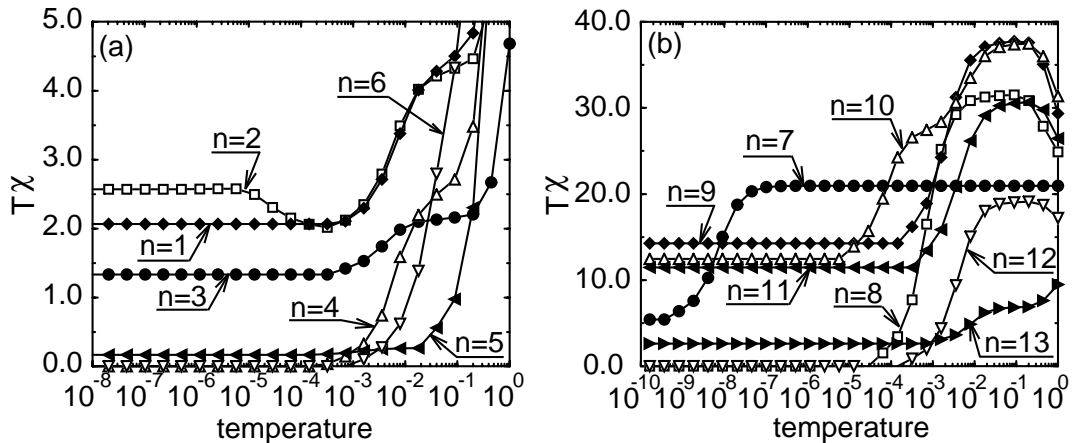


Figure 1: Magnetic susceptibility of  $f$  electron for (a)  $n=1\sim 6$  and (b)  $n=7\sim 13$ .