## (8a4)

## New High-Field Ordered State in $PrFe_4P_{12}$

T. Tayama<sup>1</sup>, J. Custers<sup>1</sup>, H. Sato<sup>1</sup>, T. Sakakibara<sup>1</sup>, H. Sugawara<sup>2</sup> and H. Sato<sup>3</sup>

<sup>1</sup>Institute for Solid State Physics, University of Tokyo, Kashiwa, 277-8581 <sup>2</sup>Faculty of Integrated Arts and Sciences, University of Tokushima, Tokushima 770-8502 <sup>3</sup>Graduate School of Science, Tokyo Metropolitan University, Hachioji, 192-0397

 $PrFe_4P_{12}$  has attracted considerable attention because of the anomalous ordered state (phase A), apparent Kondo effect and heavy fermion behavior. Recently, we showed the presence of a new ordered state (phase B) in high fields applied along the [111] direction [1]. In this talk, we discuss a possible origin of the B phase within the framework of the  $\Gamma_1-\Gamma_4$  model.

Aoki *et al.* proposed two CF schemes for  $PrFe_4P_{12}$  based on the analysis of the anisotropy in the magnetization and the electronic entropy associated with the A-phase transition, i.e.,  $\Gamma_3(0K) - \Gamma_4(18K)$  (scheme A) and  $\Gamma_1(0K) - \Gamma_4(13K)$  (scheme B) [2]. To investigate the field dependence of the two CF schemes, we carried out numerical calculations using the Hamiltonian  $H_0 = H_{\rm CF} - g_{\rm J} \mu_{\rm B} {\bf J} \cdot {\bf H}$ . Figures 1(a) and 1(b) shows the energies of the CF levels as a function of magnetic field applied along the [111] direction in schemes A and B, respectively. In scheme A, no level crossing of the two lowest levels arises, whereas a crossing occurs around 11 T in scheme B. By analogy with  $PrOs_4Sb_{12}$ , one may expect a field-induced ordered state around the level crossing point, taking into account intersite interactions. For other field orientations, no crossing of the lowest levels is found in either scheme A or B. As a result, only scheme B offers an explanation for the high-field ordered state in a limited angular range around the [111] direction. We then performed numerical calculations using scheme B and a two-sublattice mean-field Hamiltonian including quadrupolar interactions:  $H = H_0 + H_0$ . Here,  $H_{\rm Q}$  is the  $\Gamma_3$ -type quadrupolar interaction Hamiltonian, which can be given by  $H_{\rm Q} = -K[\langle O_2^0 \rangle_{\rm B(A)} O_2^0 + \langle O_2^2 \rangle_{\rm B(A)} O_2^2], O_2^0 = (3J_z^2 - \mathbf{J}^2)/2$ , and  $O_2^2 = \sqrt{3}(J_x^2 - J_y^2)/2$ . Figure 1(c) displays the calculated magnetization as functions of field at 0.5 K in the case of the coupling constant K = -0.02 K. In zero field, no ordered phase arises. Upon increasing the field, a second-order phase transition appears at 7 T, and the resulting phase diagram is given in Fig. 1(d). These calculations are qualitatively consistent with the experimental observation of phase В.

[1] T. Tayama et al.: J. Phys. Soc. Jpn. **73** (2004) 3258.

[2] Y. Aoki *et al.*: J. Phys. Soc. Jpn. **71** (2002) 2098.



Figure 1: Energy of the crystal-field levels as a function of magnetic field along the [111] axis for scheme A (a) and B (b). Magnetization curve at 0.5 K (c) and the H - T phase diagram (d) for scheme B.